

European Graduate School

**Complex Systems of  
Hadrons and Nuclei**

Copenhagen - Giessen - Helsinki - Jyväskylä



# **Neutrino interactions with nucleons and nuclei at intermediate energies**

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## **Overview:**

- Motivation & Introduction
  - Neutrino Nucleon Reactions
  - BUU Transport Model
  - Nuclear Effects in  $\nu A$  Scattering
  - Summary & Outlook
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## ■ Motivation & Introduction

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- past, current & future experiments
  - neutrino oscillations ✓
  - neutrino mass ✓
  - precision measurement of oscillation parameters ✗
  
- problems
  - uncertainties due to neutrino cross sections & nuclear effects
    - detector response
  - neutrino energy reconstruction
  - proposed experiments: MINERvA, FINESSE

➔ better understanding of nuclear effects is crucial  
for existing & future neutrino experiments

# ■ Neutrino Nucleus Scattering

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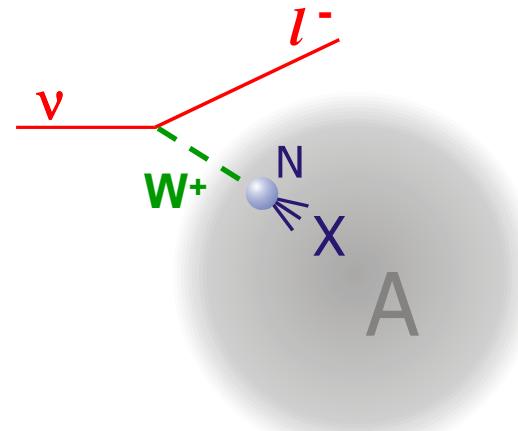
- $\nu A$  reaction is factorized using **impulse approximation**:

- $\nu_l N \rightarrow l^- X$

- with consideration of

- Fermi motion
    - Pauli blocking
    - binding energies
    - in-medium modified  $\Delta$  width

$$\Gamma \rightarrow \Gamma_{tot}^{med} = \tilde{\Gamma} + \Gamma_{coll}$$



- propagation of final state  $X$  within  
**BUU transport model** with consideration of FSI

- most general: all neutrino flavors, all nuclei, CC & NC

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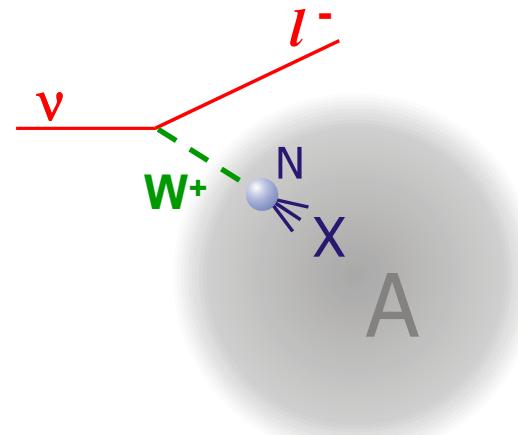
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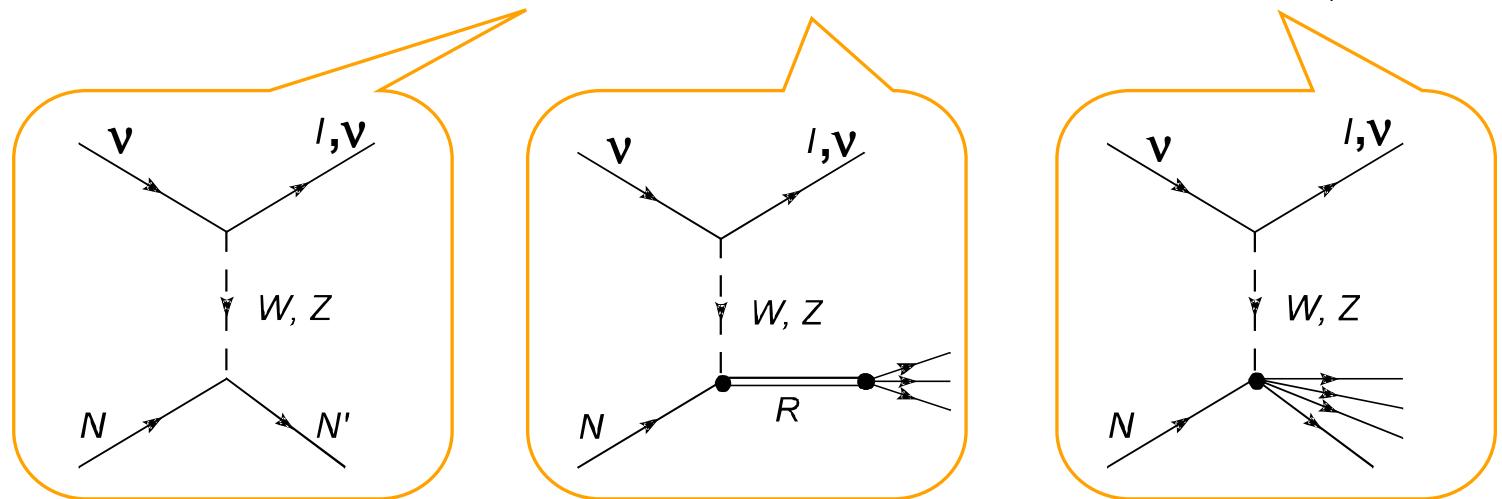


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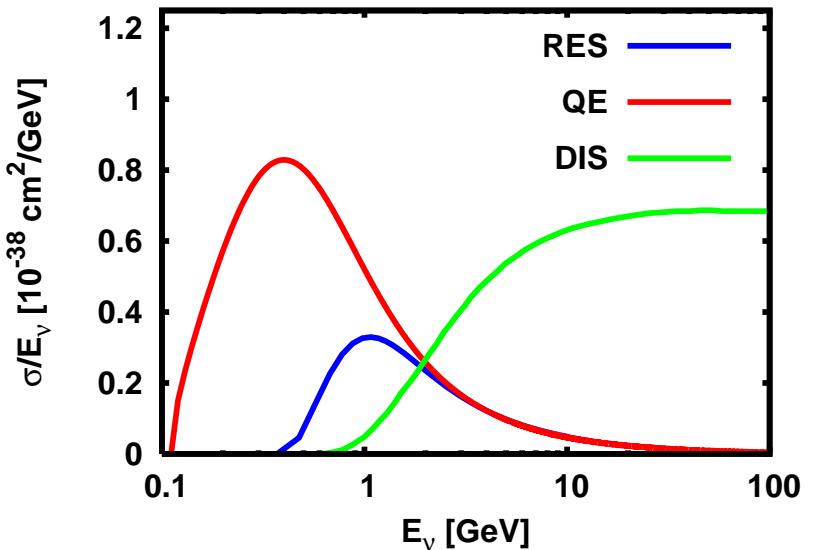
- elementary processes:  $\sigma = \sigma(QE) + \sigma(RES) + \sigma(Non-RES/DIS)$



- dominated by **QE &  $\Delta$  resonance**

CC:

$$\begin{aligned} \nu n &\rightarrow l^- p \\ \nu n &\rightarrow l^- \Delta^+ \\ \nu p &\rightarrow l^- \Delta^{++} \end{aligned}$$



# Quasielastic Scattering

- hadronic current for  $\nu_l n \rightarrow l^- p$

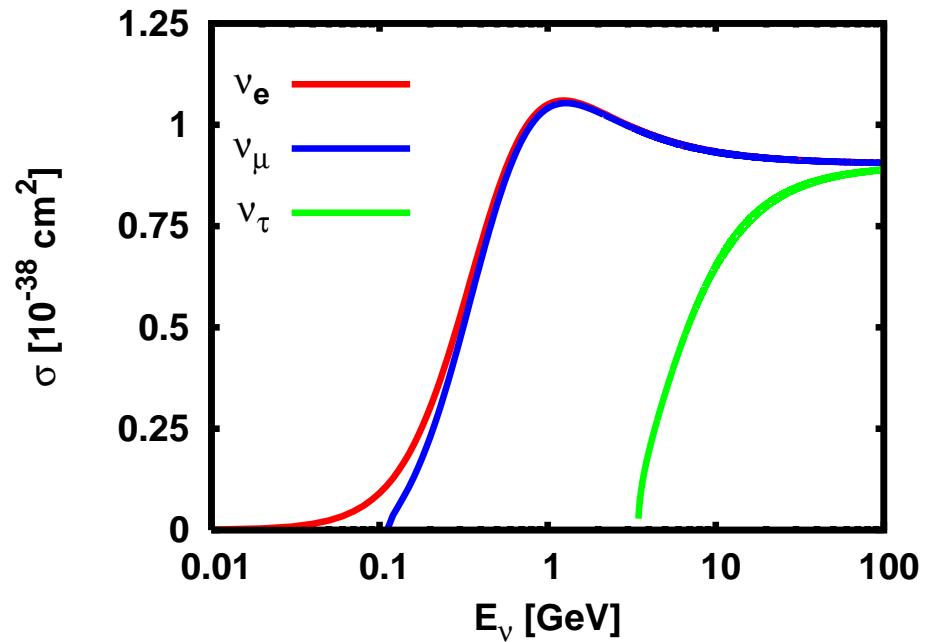
$$J^\alpha = \cos\theta_C \bar{u}_p \left( \gamma^\alpha F_1^V(Q^2) + \frac{i\sigma^{\alpha\beta} q_\beta}{2M} F_2^V(Q^2) + \gamma^\mu \gamma^5 F_A(Q^2) + \frac{q^\alpha \gamma^5}{M} F_P(Q^2) \right) u_n$$

**CVC**  
 $\downarrow$   
 $F_{1,2}^V(Q^2) = F_{1,2}^p(Q^2) - F_{1,2}^n(Q^2)$

**PCAC**  
 $\downarrow$   
 $F_P(Q^2) = \frac{2M^2}{m_\pi^2 + Q^2} F_A(Q^2)$

- BBA-2003 parametrization for  $F_{1,2}^{n,p}$  and

$$F_A(Q^2) = \frac{g_A}{\left(1 + \frac{Q^2}{M_A^2}\right)^2}$$



# ■ $\Delta$ Resonance Production

- hadronic current for  $\nu_l n \rightarrow l^- \Delta^+$

$$J_\alpha = \cos\theta_C \bar{\psi}^\beta(p') D_{\beta\alpha} u(p)$$

with the Rarita-Schwinger spinor  $\bar{\psi}^\beta(p')$  and

$$\begin{aligned} D_{\beta\alpha} = & \left( \frac{C_3^V}{M}(g_{\alpha\beta}q - q_\beta\gamma_\alpha) + \frac{C_4^V}{M^2}(g_{\alpha\beta}q \cdot p' - q_\beta p'_\alpha) + \frac{C_5^V}{M^2}(g_{\alpha\beta}q \cdot p - q_\beta p_\alpha) + g_{\alpha\beta} C_6^V \right) \gamma_5 \\ & + \frac{C_3^A}{M}(g_{\alpha\beta}q - q_\beta\gamma_\alpha) + \frac{C_4^A}{M^2}(g_{\alpha\beta}q \cdot p' - q_\beta p'_\alpha) + C_5^A g_{\alpha\beta} + \frac{C_6^A}{M^2} q_\beta q_\alpha \end{aligned}$$

**CVC &  $M_{1+}$  dominance**

$$\begin{gathered} \downarrow \\ C_4^V \sim C_3^V \quad C_5^V = 0 \quad C_6^V = 0 \\ C_3^V \rightarrow eN \end{gathered}$$

**PCAC**

$$\downarrow \quad C_6^A \sim C_5^A$$

**parametrization**

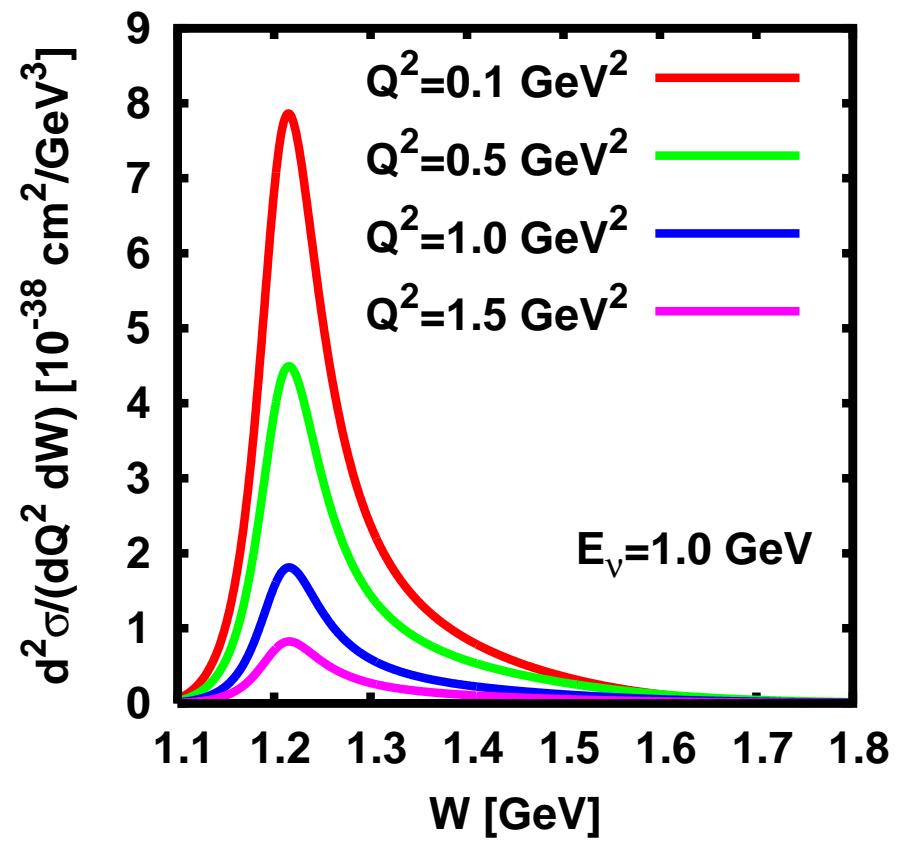
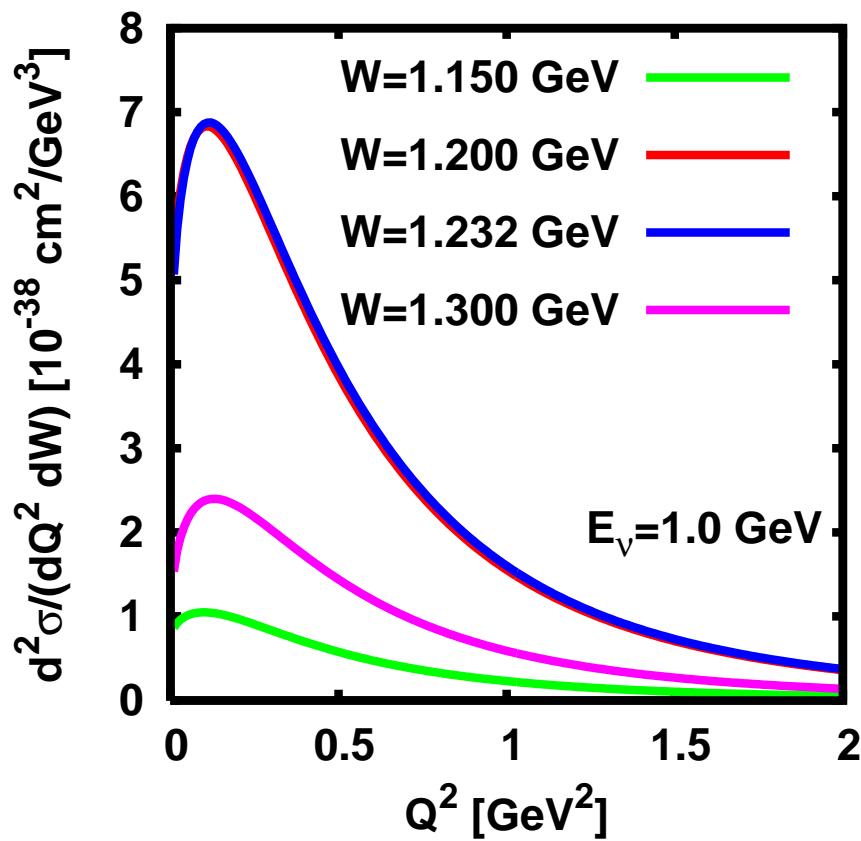
$$\downarrow \quad \begin{matrix} C_3^V \\ C_5^A \quad C_4^A \quad C_3^A \end{matrix}$$

- $\Delta$  width: **p-wave**  $\Gamma \sim q_{CM}^3$

# ■ $\Delta$ Resonance Production Cross Section

- double differential cross section

$$\frac{d^2\sigma}{dQ^2 dW} \text{ for } \nu \mu p \rightarrow \mu^- \Delta^{++}$$



# ■ Neutrino Nucleus Scattering

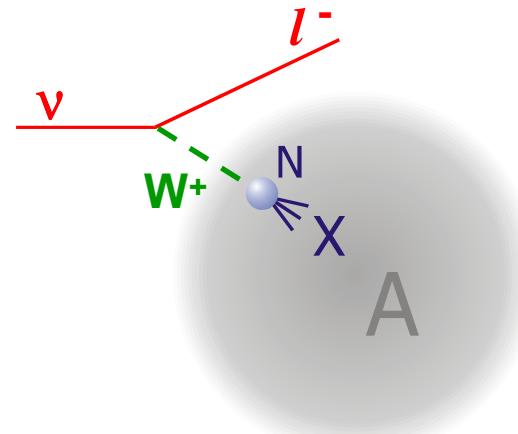
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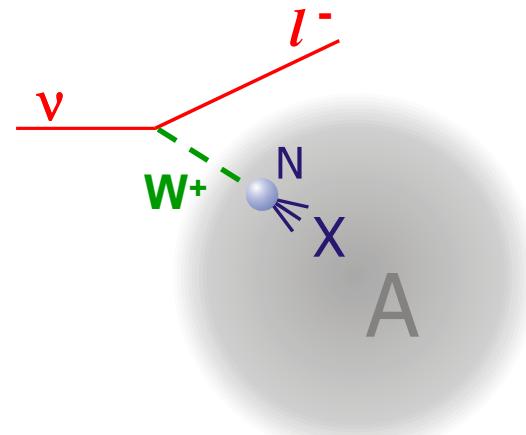
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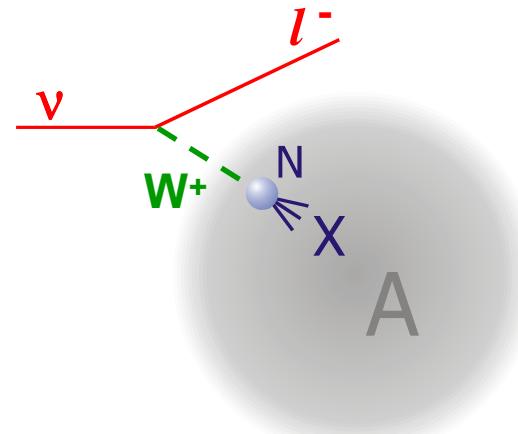
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## ■ BUU Transport Model

- description of heavy ion collisions,  $e A$ ,  $\gamma A$  and  $\nu A$  reactions with **one** code
- compared experimental data, in particular for  $\gamma^{(*)} A$
- coupled channel semiclassical transport model
  - Boltzmann-**Uehling-Uhlenbeck** equation  
for each particle species  $i$  ( $i = N, R, \pi, \rho, K, \dots$ ):

$$\frac{df_i}{dt} = \left( \partial_t + (\nabla_{\vec{p}} H) \nabla_{\vec{r}} - (\nabla_{\vec{r}} H) \nabla_{\vec{p}} \right) f_i(\vec{r}, \vec{p}, t) = I_{coll} [f_1, \dots, f_i, \dots, f_M]$$

Hamiltonian:  $H = \sqrt{(\mu + U_s)^2 + \vec{p}^2}$        $f_i$  : phase space density



**mean field for baryons**

Skyrme type with momentum dependence

- set of BUU equations coupled via  $I_{coll}$  and mean field

# ■ BUU Transport Model – Collision Term

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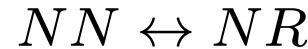
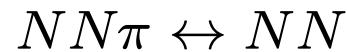
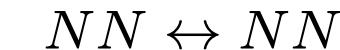
## ■ **collision integral** : 2 particle collisions:

- elastic and inelastic scattering (coupled channels)
- Pauli blocking for fermions

**FSI**    

- absorption
- charge exchange
- redistribution of energy
- production of new particles

## ■ most important scattering processes:



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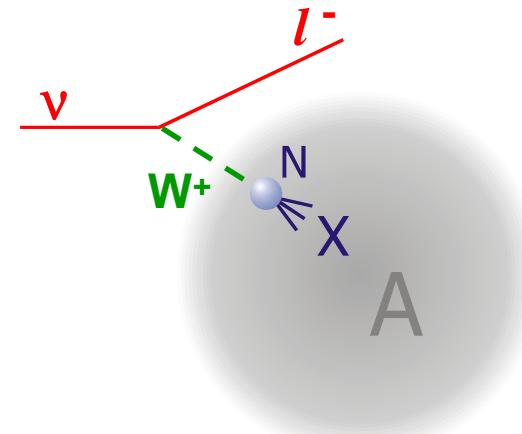
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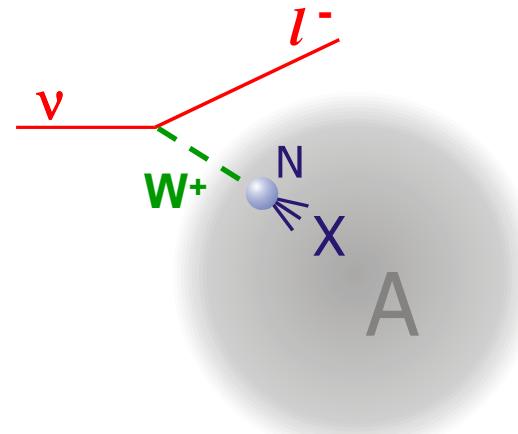
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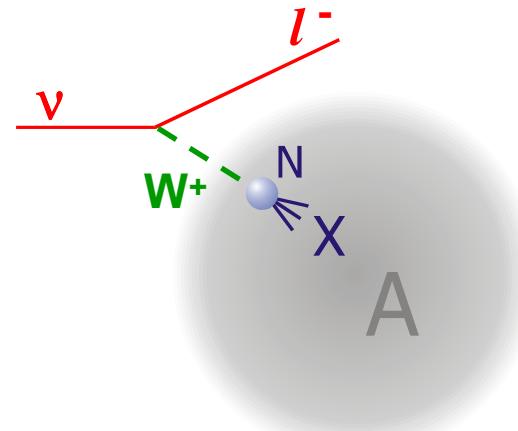
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- propagation of final state  $X$  within the **BUU transport model** taking into account **FSI**

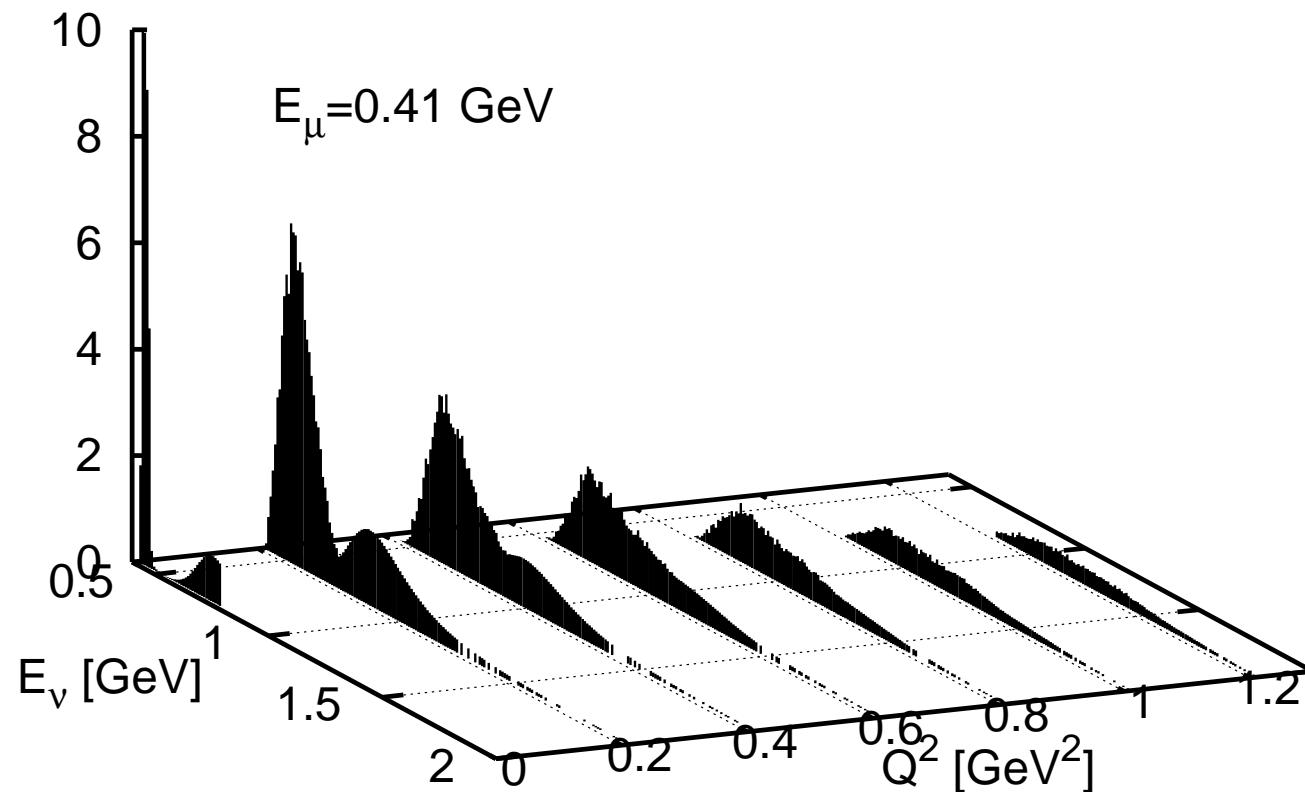
- example:  $\nu_\mu {}^{56}\text{Fe} \rightarrow \mu^- X$  at  $E_\nu = 0.4 - 2 \text{ GeV}$

## ■ Inclusive Cross Section $\nu_\mu {}^{56}\text{Fe} \rightarrow \mu^- X$

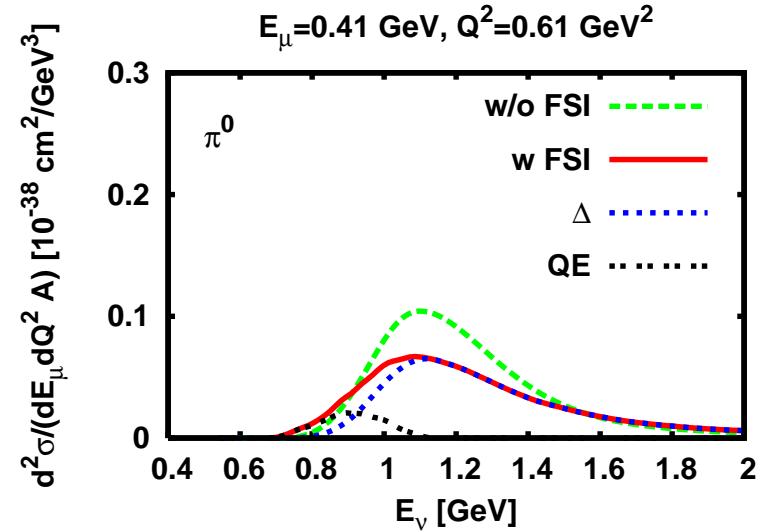
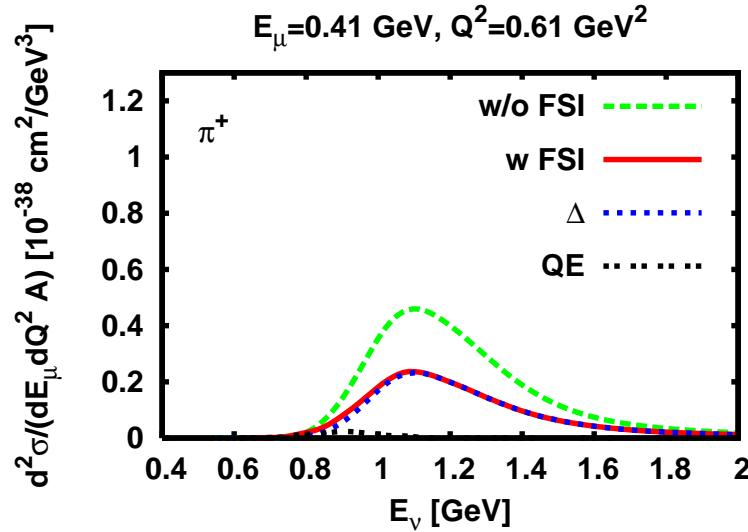
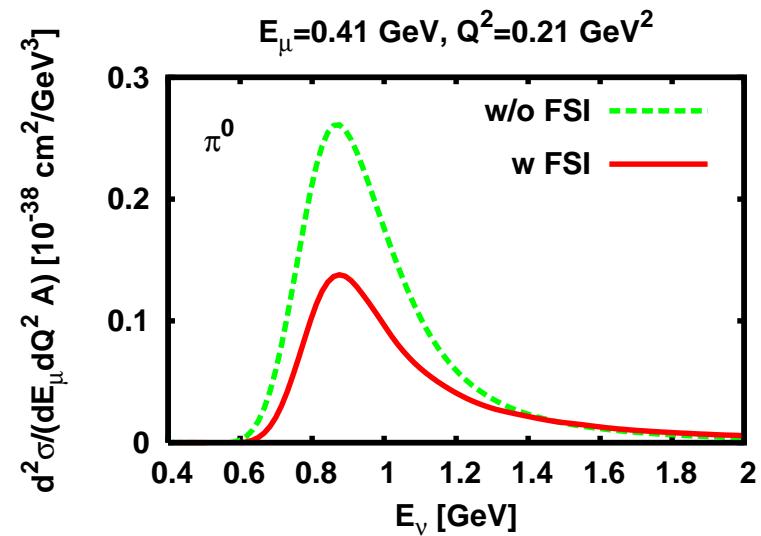
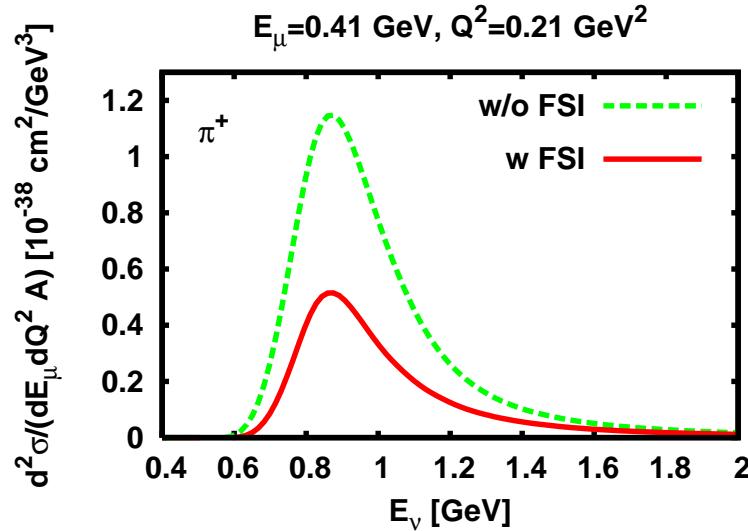
- double differential cross sections per nucleon:  $\frac{d^2\sigma}{dE_\mu dQ^2 A}$

- inclusive cross section:

$$d^2\sigma/(dE_\mu dQ^2 A) [10^{-38} \text{ cm}^2/\text{GeV}^3]$$

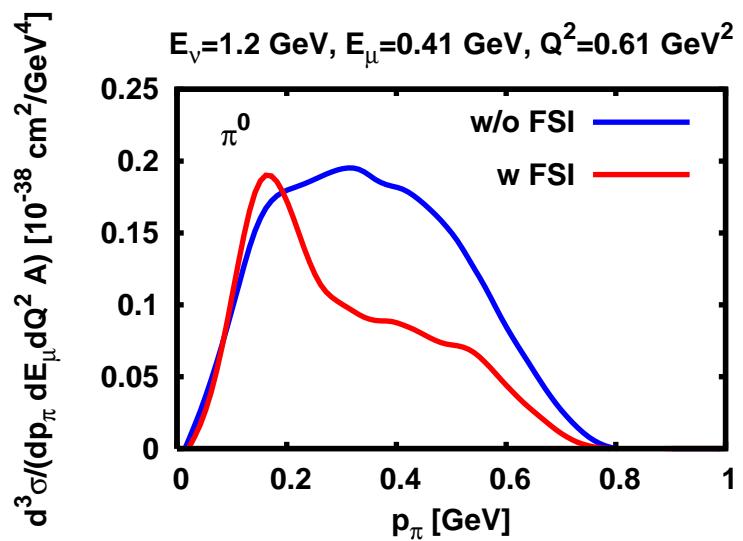
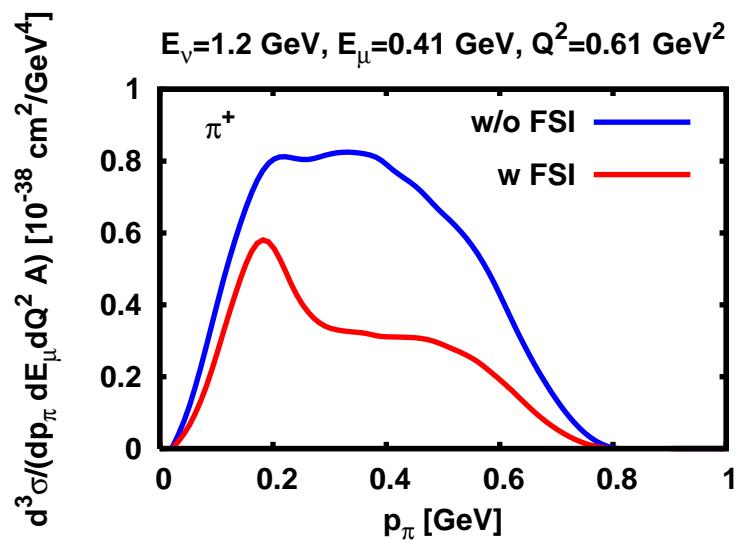
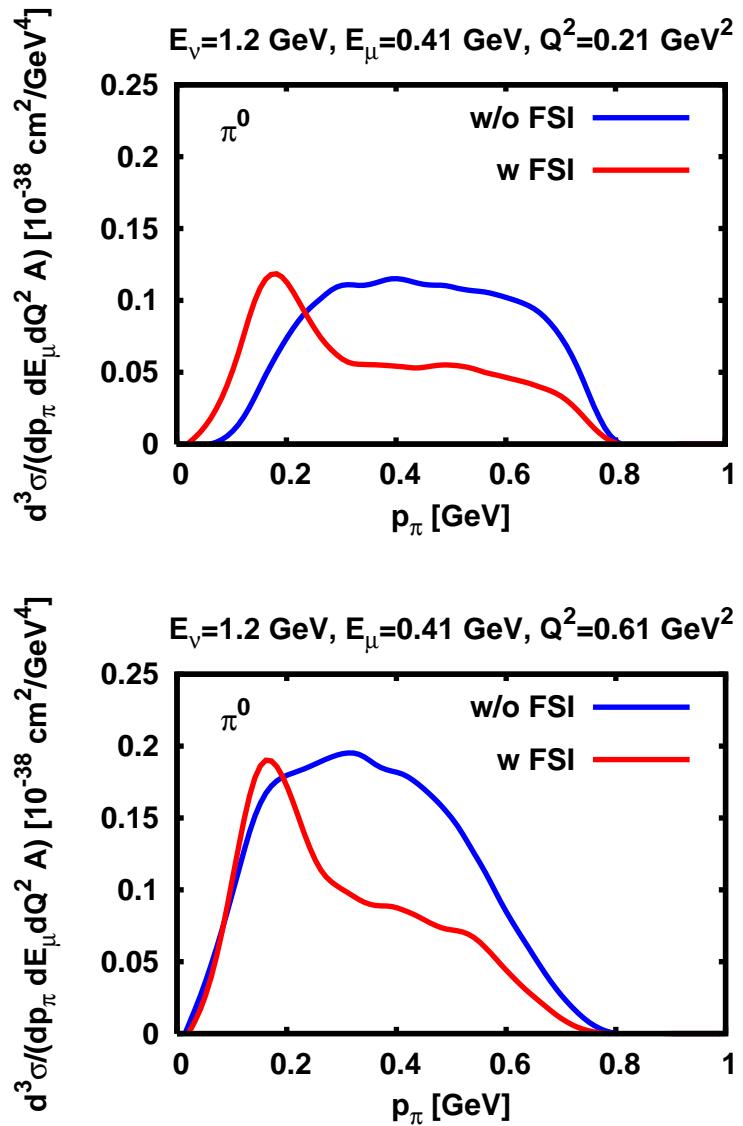
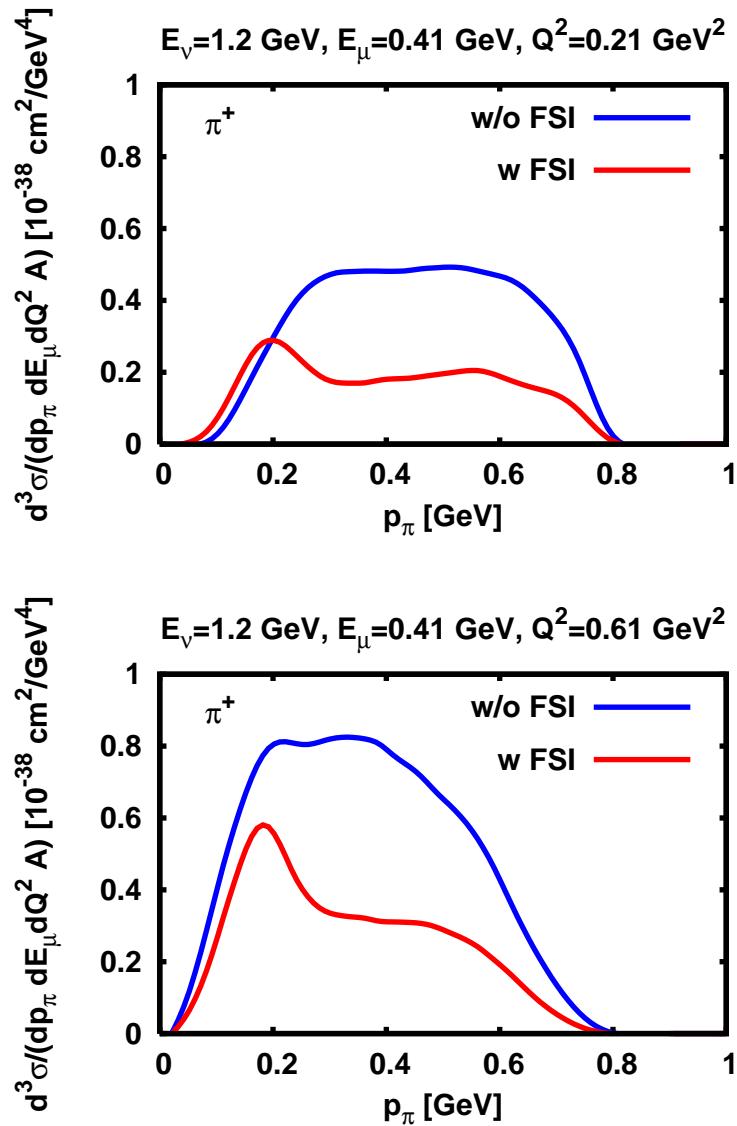


# Pion Production $\nu_\mu {}^{56}\text{Fe} \rightarrow \mu^- \pi X$



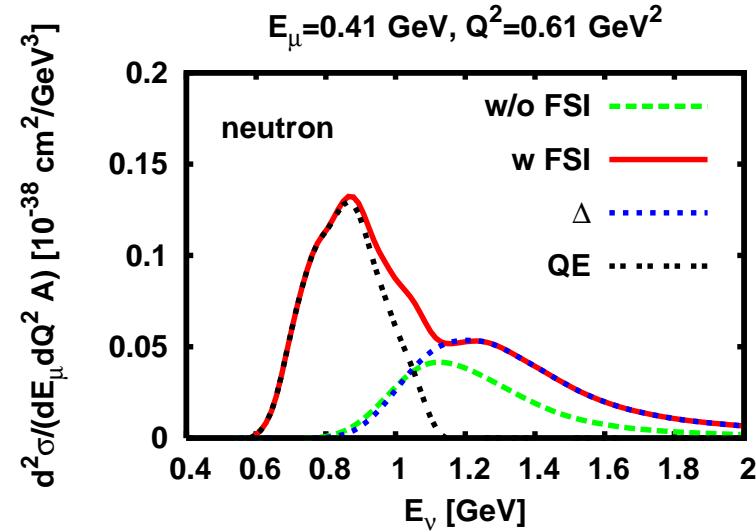
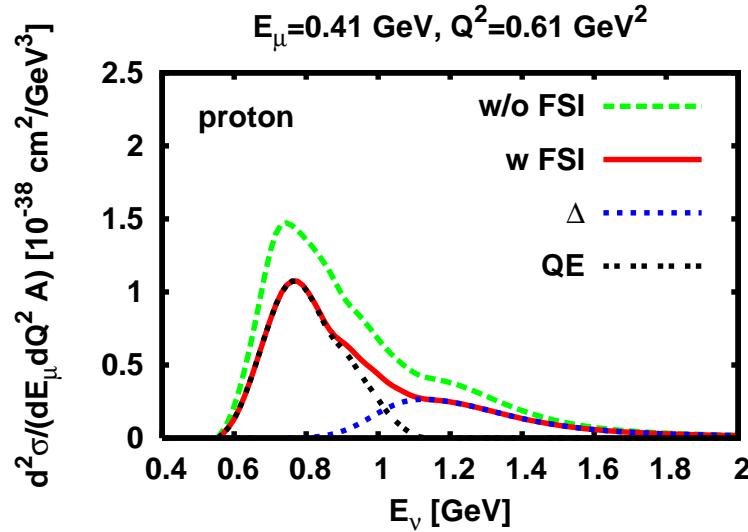
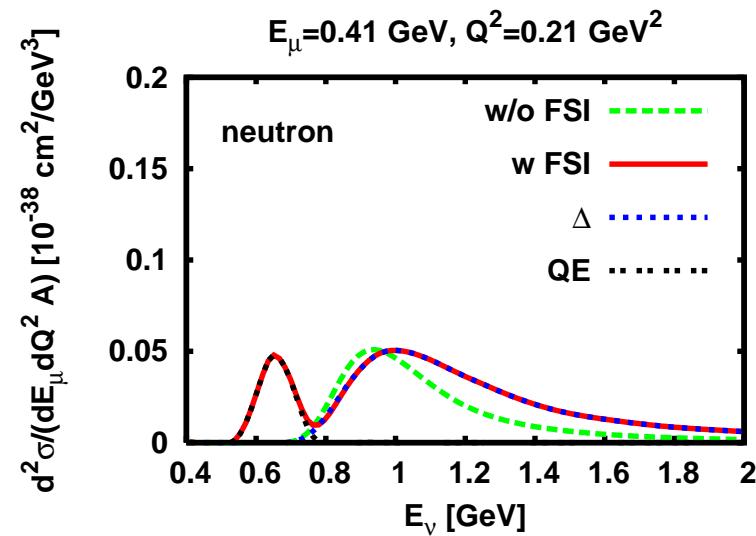
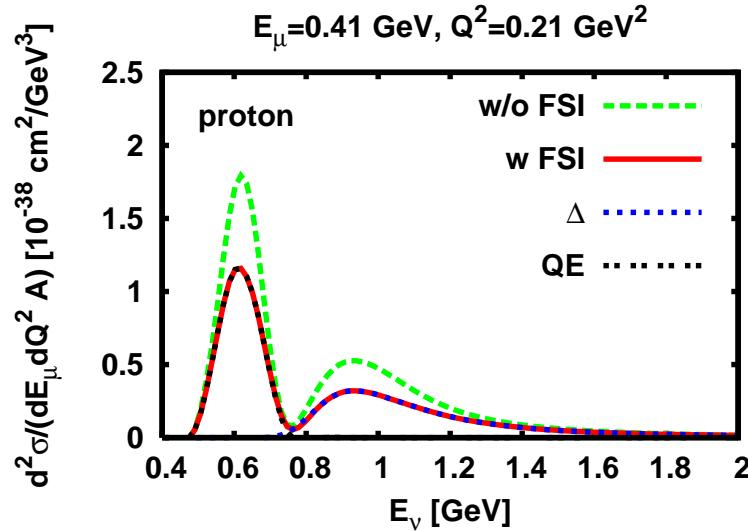
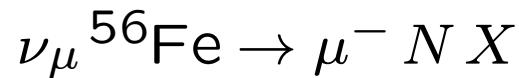


# Pion Momentum Distribution $\nu_\mu {}^{56}\text{Fe} \rightarrow \mu^- \pi X$





# Nucleon Knockout



# ■ Summary & Outlook

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## ■ neutrino nucleus scattering

- impulse approximation
- $\nu N$  followed by FSI

## ■ neutrino nucleon reactions

- quasielastic scattering &  $\Delta$  production
- vector form factors extracted from electron scattering

## ■ BUU model

- extended to  $\nu A$
- important in-medium effects are taken into account

## ■ nuclear effects in $\nu A$ scattering

- inclusive scattering, pion production & nucleon knockout
- in-medium effects, in particular FSI, are not negligible